

Strong uniaxial pressure dependencies evidencing spin-lattice coupling and spin fluctuations in $\text{Cr}_2\text{Ge}_2\text{Te}_6$

S. Spachmann¹, S. Selzer², B. Büchner^{2,3}, S. Aswartham², and R. Klingeler^{1,*}

¹Kirchhoff Institute for Physics, Heidelberg University, INF 227, 69120 Heidelberg, Germany

²Leibniz Institute for Solid State and Materials Research (IFW), Helmholtzstraße 20, 01069 Dresden, Germany

³Würzburg-Dresden Cluster of Excellence *ct.qmat*



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Single crystals of $\text{Cr}_2\text{Ge}_2\text{Te}_6$ were studied by high-resolution capacitance dilatometry to obtain in-plane ($B \parallel ab$) and out-of-plane ($B \parallel c$) thermal expansion and magnetostriction at temperatures between 2 and 300 K and in magnetic fields up to 15 T. The anomalies in both response functions lead to the “magnetoelastic” phase diagrams and separate the paramagnetic (PM), ferromagnetic low-temperature/low-field (LTF) and aligned ferromagnetic (FM) phases. Different signs of magnetostriction anomalies as well as the evolution of thermal expansion anomalies at small fields $B \parallel ab$ of different magnetic-field dependence clearly supports the scenario of an intermediate region separating PM and LTF phases in finite external in-plane magnetic fields and implies a triple point in the magnetic phase diagram. Simulations of the magnetostriction using the Stoner-Wohlfarth model for uniaxial anisotropy demonstrate that the observed quadratic-in-field behavior in the LTF phase is in line with a rotation of the spins from the preferred c direction into the ab plane. Both the LTF and the PM phase close to T_C exhibit very strong pressure dependencies of the magnetization, $\partial \ln M_{ab} / \partial p_{ab}$, of several hundred %/GPa and the transition from the LTF to the FM phase strongly depends on p_{ab} ($\sim 280\%$ /GPa), indicating a strong decrease in the uniaxial anisotropy under applied in-plane pressure. Our data clearly demonstrate the relevance of critical fluctuations and magnetoelastic coupling in $\text{Cr}_2\text{Ge}_2\text{Te}_6$.

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I. INTRODUCTION

Magnetic anisotropy is one of the key features of quasi-two-dimensional (quasi-2D) magnetic van der Waals (vdW) materials [1]. Purposefully engineering this anisotropy, e.g., by strain [2–5], pressure [6–9], doping [10,11], or the combination with other materials in heterostructures [12–14], is a crucial ingredient in the hunt for room-temperature 2D ferromagnets on one end of the spectrum between applied and fundamental research and for (Kitaev) quantum spin liquids (QSLs) on the other end [15,16]. Achieving room-temperature 2D ferromagnetism in vdW materials is desirable for a large number of applications, among them spintronic, magnonic, and spin-orbitronic devices [1]. Identifying and engineering (Kitaev) QSLs on the other hand, on top of the significance for fundamental physics, comes with the promise of topologically protected quantum computing [17] due to presence of Majorana fermions [18]. In the quests to engineer the properties and especially the anisotropy of magnetic quasi-2D materials, the interlayer and intralayer spacings play a significant role. Only small changes in these spacings can lead to drastic responses in magnetic and electronic properties [19]. These drastic responses originate from the delicate interplay of charge, spin, and lattice degrees of freedom in correlated electron systems.

The transition-metal trichalcogenide $\text{Cr}_2\text{Ge}_2\text{Te}_6$ is one prominent example of the currently highly investigated

layered magnetic vdW materials. Its layers, stacked along the c axis, consist of a honeycomb network of edge-sharing CrTe_6 octahedra. $\text{Cr}_2\text{Ge}_2\text{Te}_6$ crystallizes in the trigonal space-group $R\bar{3}$ (No. 148) and becomes ferromagnetic below $T_C \approx 65$ K in its bulk form. This ferromagnetism is preserved when the material is thinned to a bilayer where long-range magnetic order is observed below about 30 K [20].

While there is consensus that $\text{Cr}_2\text{Ge}_2\text{Te}_6$ exhibits uniaxial anisotropy at ambient pressure, conflicting claims about a possible change from uniaxial to easy-plane anisotropy under hydrostatic pressure above 1 GPa have been made [21,22]. Importantly, strong spin-lattice coupling has been proven by Raman scattering at ambient and nonambient pressures [23,24]. Furthermore, zero-field thermal expansion measurements of $\text{Cr}_2\text{Ge}_2\text{Te}_6$ have revealed sharp anomalies both along the in-plane and out-of-plane directions associated with the evolution of long-range magnetic order at T_C , evidencing the presence of significant magnetoelastic coupling [25].

In this paper, we investigate the effect of magnetic fields on the structural and magnetic properties of $\text{Cr}_2\text{Ge}_2\text{Te}_6$. In particular, we investigate the effects on the in-plane and out-of-plane lattice parameters with a special focus on the low-temperature and low-field (LTF) phase, which is present when a small magnetic field is applied along the in-plane direction. The magnetostriction and magnetization below T_C for $B \parallel ab$ are modeled using the Stoner-Wohlfarth model. We establish the magnetoelastic phase diagram and derive and discuss the uniaxial pressure dependencies of the critical field and of the magnetization.

*klingeler@kip.uni-heidelberg.de

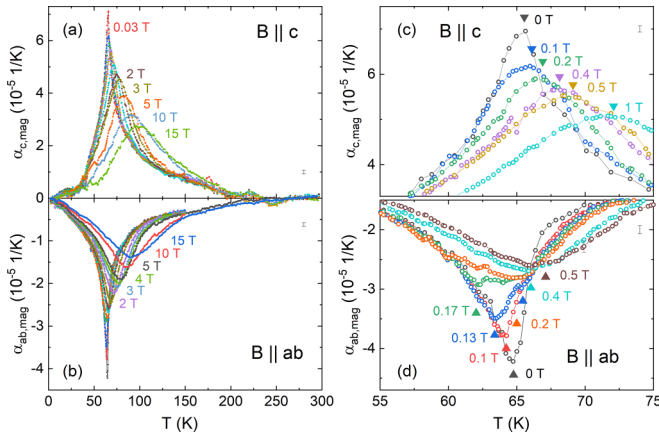


FIG. 1. Magnetic contributions to the thermal expansion coefficient in external fields between 0 and 15 T (a) and (b) and between 0 and 1 T (c) and (d) for $B \parallel c$ and $B \parallel ab$, respectively. Triangles in (b) and (c) mark the peak positions (see also Fig. S2 in the Supplemental Material [30]).

II. EXPERIMENTAL DETAILS

Single crystals of $\text{Cr}_2\text{Ge}_2\text{Te}_6$ have been grown by the self-flux technique and were characterized in detail as reported in Refs. [26,27]. The crystals measured in this paper have a very thin cuboid shape with a length of 0.255 mm along the c axis for one crystal, as well as in-plane dimensions of $1.3 \times 2.0 \text{ mm}^2$ for a second crystal used for in-plane measurements. High-resolution dilatometry measurements were performed by means of a three-terminal high-resolution capacitance dilatometer from Kuechler Innovative Measurement Technology in a homebuilt setup [28,29]. The capacitance read-off was facilitated by Andeen-Hagerling's AH 2550A Ultra-Precision 1-kHz capacitance bridge. With the dilatometer, the uniaxial relative length changes $\Delta L_i(T)/L_i$ and the linear thermal expansion coefficients $\alpha_i = 1/L_i dL_i(T)/dT$ both along the c axis and along the in-plane direction, i.e., $\parallel ab$, were studied at temperatures between 2 and 300 K in zero field and in magnetic fields up to 15 T. Magnetic fields were applied along the direction of the measured length changes $i = c, \parallel ab$. In addition, the field-induced length changes $\Delta L_i(B_i)$ were measured at various fixed temperatures between 2 and 204 K in magnetic fields up to 15 T. Magnetization measurements were performed in a MPMS3 SQUID-Magnetometer by Quantum Design using the vibrating sample magnetometer option.

III. RESULTS

A. Thermal expansion and magnetostriction

Figure 1 shows the effect of magnetic field applied along the c axis (ab plane) on the thermal expansion coefficient α_c (α_{ab}). The same phonon background as in Ref. [25] was subtracted to obtain the nonphononic contributions $\alpha_{i,\text{mag}}$ shown in the figure. At zero field, pronounced anomalies in the thermal expansion coefficients imply significant magnetoelastic coupling and confirm the presence of strong spin fluctuations at T_C . As a general trend for both measurement directions, the anomalies in $\alpha_{i,\text{mag}}$ around T_C broaden and shift to higher tem-

peratures as the magnetic field is increased. At small magnetic fields $B \parallel ab$ below 0.2 T, however, an initial suppression of the peak position is observed. The suppression of the anomaly associated with the onset of the LTF phase implies the presence of a competing ferromagnetic phase, which is stabilized in magnetic fields with respect to the LTF phase (whereas the crossover temperature into the PM phase increases as expected for a ferromagnetic transition). Indeed, the data at 0.13 T show the presence of an additional small feature in the high-temperature shoulder of the anomaly (Fig. 1(d) and Fig. S2 in the Supplemental Material [30]). The lower-temperature anomaly confirms the initial suppression of the dominating peak to 63.3 K. It is accompanied by a weak feature at 65.4 K (i.e., slightly above T_C). At 0.17 T, the whole feature further broadens and extends to higher temperature whereas only a small additional feature remains at 62.3 K. As the field is increased above 0.17 T, the broad anomaly continues to shift to higher temperatures whereas the small one cannot be distinguished anymore. The different field dependencies and different signs of magnetostriction anomalies [see Fig. 5(b)] clearly imply the presence of an intermediate phase separating the LTF and PM phases as suggested from the analysis of specific heat and magnetization data in Ref. [27]. This is further confirmed by the observed evolution of the thermal expansion anomalies in small fields $B \parallel ab$ which shows that the FM phase extends to small fields. In contrast to Ref. [27], the data at hand indicate that the two phase boundaries merge to a single point slightly above or at $B = 0$ T. The merged phase boundaries might be characterized by tricritical behavior which has been found by Liu *et al.* from their analysis of the critical behavior around T_C [31].

In-plane magnetostriction measurements with $B \parallel ab$ reveal a behavior which clearly implies the presence of two distinct phases below T_C [Figs. 2(a) and 2(b)]. At low fields and temperatures the relative length changes $\Delta L_{ab}(B)$ exhibit a roughly quadratic-in-field decrease $\Delta L_{ab}(B) \propto B^2$ before transitioning to a constant value, i.e., vanishing magnetostriction. We note that the prefactor of the B^2 dependence for the quadratic-in-field magnetostriction very well follows the temperature dependence of the magnetic order parameter given by the refined magnetic moment of Cr^{3+} ions [32] (see Fig. S3 in the Supplemental Material [30]). Increasing the temperature towards T_C decreases the range of the quadratic-in-field behavior and the magnetostriction coefficient above the transition assumes a positive value. The jump in λ_{ab} suggests that the transition between the two regimes may be interpreted as a second-order phase transition.

In-plane measurements of the isothermal magnetization confirm the observed behavior. The magnetic susceptibility $\partial M_{ab}/\partial B$ exhibits a jump at the same fields as λ_{ab} [Fig. 2(c)]. In contrast, magnetostriction and magnetization measurements with $B \parallel c$ below T_C show no phase transition (see the Supplemental Material [30], Fig. S4).

B. Simulations using the Stoner-Wohlfarth model

In order to model the low-temperature magnetization and magnetostriction for $B \parallel ab$ the Stoner-Wohlfarth (SW) model [33] was applied. The normalized energy density in this model—with ϕ being the angle between the magnetic

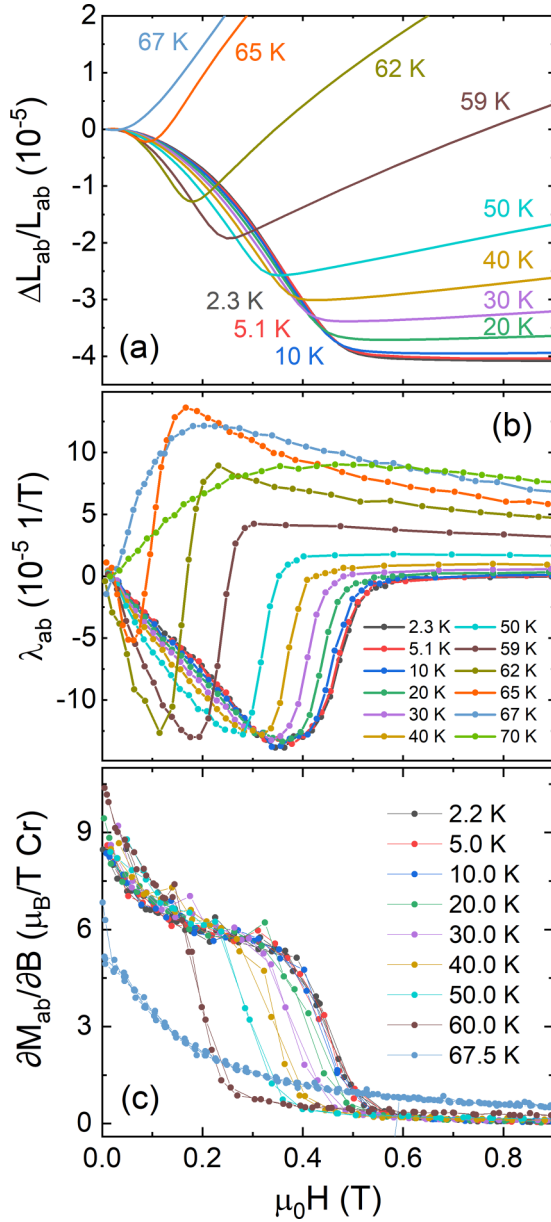


FIG. 2. Relative length changes (a), magnetostriction coefficient (b), and magnetic susceptibility (c) at low fields for $B \parallel ab$ and temperatures from 2 to 70 K. Only up-sweep data are shown in (a) and (b), whereas up- and down-sweeps are shown in (c).

moment, expressed by the saturation magnetization M_S , and external magnetic-field H , and θ being the angle between the easy axis and H —is given by

$$\eta = \frac{E}{2K_U V} = \frac{1}{4} - \frac{1}{4} \cos 2(\phi - \theta) - h \cos \phi, \quad (1)$$

where the first two terms describe the anisotropy energy density and the third term the Zeeman energy density. The normalized field is given by $h = \mu_0 M_S H / (2K_U)$, where μ_0 is the vacuum permeability, M_S is the saturation magnetization, and K_U is the energy density characterizing the uniaxial anisotropy along c . Since the field was applied on the ab plane, i.e., $\perp c$, we have $\theta = \pi/2$. η was minimized with respect to ϕ to obtain the direction of the magnetic moment at a given field.

Due to the uniaxial symmetry of the system two values ϕ_i minimize η for $H < H_{\text{sat}}$. The magnetization along the applied field direction is then obtained as

$$M(h, T) = \sum_{i=1}^2 M_S(T) \cos[\phi_i(h)] \sin \theta_i / 2 \quad (2)$$

$$= \sum_{i=1}^2 M_S(T) \cos[\phi_i(h)] / 2. \quad (3)$$

For $H \geq H_{\text{sat}}$, the magnetic moment is aligned along the field, i.e., $\phi = 0$ and $M(h, T) = M_S(T)$. $M_S(T)$ is extracted from the magnetization at the intersection of linear fits to the $M(B)$ data at the respective temperature in the field range from 0 to 0.2 T and 3 to 7 T (at 60 K: 0 to 0.15 T and 4.5 to 7 T). Analogous to the magnetization, the magnetostriction is given by

$$\Delta L(h, T) = \frac{K_{U,\text{eff}}(T)}{K_U} \lambda_{\text{sat}} \sum_{i=1}^2 \{\cos[\phi_i(h)] \sin \theta_i\}^2 / 2 \quad (4)$$

$$= \frac{K_{U,\text{eff}}(T)}{K_U} \lambda_{\text{sat}} \sum_{i=1}^2 \{\cos[\phi_i(h)]\}^2 / 2, \quad (5)$$

where λ_{sat} is the saturation magnetostriction along the measurement direction and $K_{U,\text{eff}}$ is an effective anisotropy parameter. This effective anisotropy parameter is calculated according to the theory by Callen and Callen [34] as

$$\frac{K_{U,\text{eff}}(T)}{K_U} = \left[\frac{M_S(T)}{M_S} \right]^{l(l+1)/2}. \quad (6)$$

In the following, we approximate $K_U \approx K_{U,\text{eff}}$ (2.2 K) and $M_S \approx M_S(2.2 \text{ K})$.

Applying the model as described above, the best agreement between magnetostriction results at 2.2 K (solid line) and the simulation (open circles) is obtained for $K_U = 46.35 \text{ kJ/m}^3$ and $\lambda_{\text{sat}} = -4.08 \times 10^{-5}$ [Fig. 3(a)]. Small deviations between experimental results and simulations are visible at low fields where the down-sweep data is closer to the simulated values. Also, a smooth transition is observed in the experiments signaling the full alignment of the magnetic moments with the magnetic field at roughly 0.5 T in contrast to a kink in the numerical results. Overall the data are described very well by the simulation. The same holds for the magnetization at low temperatures [Fig. 3(c)].

Increasing the temperature towards T_C , the experimental data (solid lines) can be described reasonably well using $l = 1.83$ for the calculations (open circles), which is close to $l = 2$ [$n = l(l+1)/2 = 3$] expected for uniaxial anisotropy. Up to 20 K and 1 T these parameters describe the experimental results well, however, above this temperature and field two types of strong deviations can be distinguished: (1) the experimental magnetostriction is larger at low fields than the simulated one, and (2) a change from zero to positive magnetostriction above the saturation field becomes visible which gets larger in magnitude as the temperature approaches T_C [see Fig. 3(b)]. The latter effect is most likely caused by a disturbance of the spin alignment due to phonons and spin-phonon coupling as the temperature is increased. This explanation is supported by the experimental magnetization, which is significantly lower than

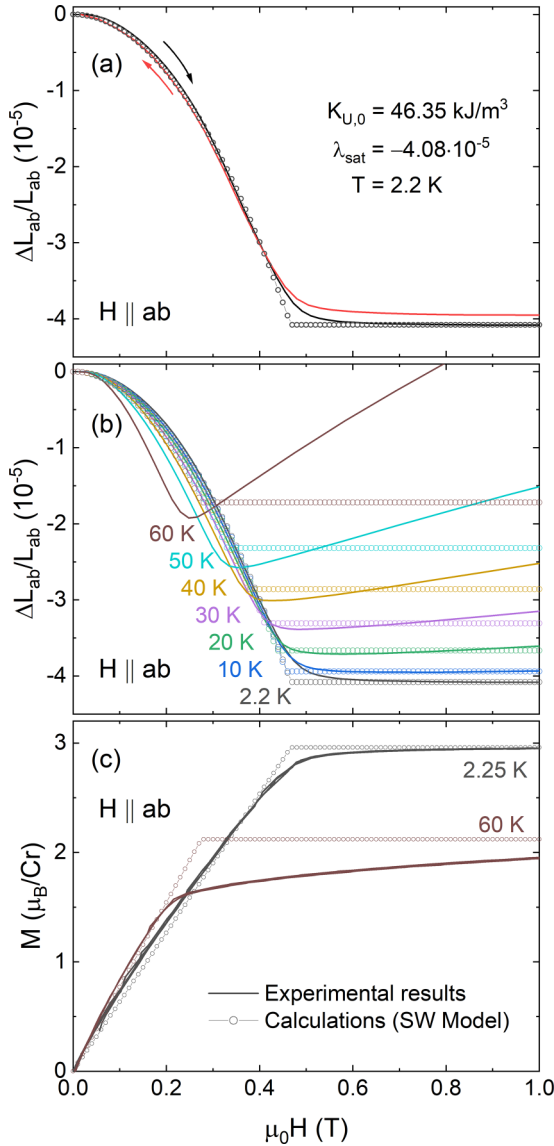


FIG. 3. Comparison of experimental magnetostriction data (a) and (b) and magnetization data (c) for $H \parallel ab$ (solid lines) with calculations using the Stoner-Wohlfarth model (open circles). Only up-sweep data are shown in (b), whereas both up- and down-sweep data are shown in (a) and (c).

the one calculated within the SW model [Fig. 3(c)], as well as the previous observation of a significant spin-phonon coupling in $\text{Cr}_2\text{Ge}_2\text{Te}_6$ [23]. Note that using an effective anisotropy parameter $K_{U,\text{eff}}$ does not fully cover the effects induced by raising the temperature.

IV. DISCUSSION

A. Low-temperature and low-field phase

The simulations described above show that the magnetostriction and magnetization of the LTF phase in $\text{Cr}_2\text{Ge}_2\text{Te}_6$ can largely be described by the SW model at low temperatures. This result suggests that a magnetic field applied along the in-plane direction in the LTF phase leads to a rotation of the magnetic moments from the anisotropy-preferred c

direction towards the magnetic field until saturation, i.e., full alignment, is reached. Such a rotation of the magnetic moments in the LTF phase of $\text{Cr}_2\text{Ge}_2\text{Te}_6$ was first described by S. Selzer *et al.* [27]. We also note that the SW model was also used to model ferromagnetic resonance data [26,35]. These previous reports obtained anisotropy parameters $K_U = 47(1)$ [27], 48 [26], and 39.5 kJ/m^3 [35]. The first two are in perfect agreement with $K_U = 46.35 \text{ kJ/m}^3$ found from our simulations.

B. Uniaxial pressure dependence of the critical field

For continuous phase transitions the uniaxial pressure dependence of the critical field is given as

$$\frac{\partial B_{\text{cr}}}{\partial p_{ab}} = \frac{\Delta\lambda}{\Delta\left(\frac{\partial M}{\partial B}\right)}, \quad (7)$$

where the two quantities in the numerator and denominator on the right-hand side are jumps in the magnetostriction coefficient and in the magnetic susceptibility at the phase boundary, respectively. For $\text{Cr}_2\text{Ge}_2\text{Te}_6$, these jumps amount to $\Delta\lambda_{ab} = 1.9(3) \times 10^{-4}/T$ and $\Delta(\partial M_{ab}/\partial B) = -4.6(5) \mu_B T^{-1} \text{Cr}^{-1}$ at 2.2 K, which, using the molar volume of $V_m = 1.67 \times 10^{-4} \text{ m}^3/\text{mol}$, leads to $\partial B_{\text{cr}}/\partial p_{ab} = -1.25(18) \text{ T/GPa}$ ($\simeq -280 \text{ \%/GPa}$). Up to T_C this value remains constant within error bars (see the Supplemental Material [30], Table 1). Our data, hence, imply that B_{cr} is extremely sensitive to in-plane pressure. Considering the rather small critical fields of less than 0.5 T, this result implies that the LTF phase can be fully suppressed at small uniaxial pressures so that the field-induced high-field phase will appear already at $B = 0 \text{ T}$.

This behavior suggests a strong decrease in the uniaxial anisotropy parameter under pressure $p \parallel ab$. A comparable decrease of the magnetic anisotropy energy by compressive biaxial (in-plane) strain has been calculated for a heterostructure of $\text{Cr}_2\text{Ge}_2\text{Te}_6$ [36] and for a similar van der Waals material Fe_3GeTe_2 [37]. We also note that switching from uniaxial to easy-plane anisotropy has been reported in $\text{Cr}_2\text{Ge}_2\text{Te}_6$ for hydrostatic pressures larger than 1 GPa [21], which was, however, not confirmed at such low pressures by another study [22]. In contrast to uniaxial pressure applied on the ab plane, uniaxial pressure $p \parallel c$ has been shown to strongly stabilize the FM low-temperature phase at the cost of the high-temperature paramagnetic one [25], suggesting an enhancement of the uniaxial anisotropy along c .

C. Uniaxial pressure dependence of the magnetization

By exploiting Maxwell relations, the experimentally obtained magnetostriction coefficients λ_i can be used to obtain the uniaxial pressure dependencies of the magnetization, i.e.,

$$\lambda_{ij} = \frac{1}{L_i} \left(\frac{\partial L_i}{\partial B_j} \right) = \frac{1}{L_i} \left(\frac{\partial^2 G}{\partial p_i \partial B_j} \right) = - \frac{\partial M_i}{\partial p_j}, \quad (8)$$

where G is the Gibbs potential such that the relative pressure dependence of the magnetization can be calculated as

$$\frac{\lambda_{ij}}{M_i} = - \frac{\partial \ln(M_i)}{\partial p_j}. \quad (9)$$

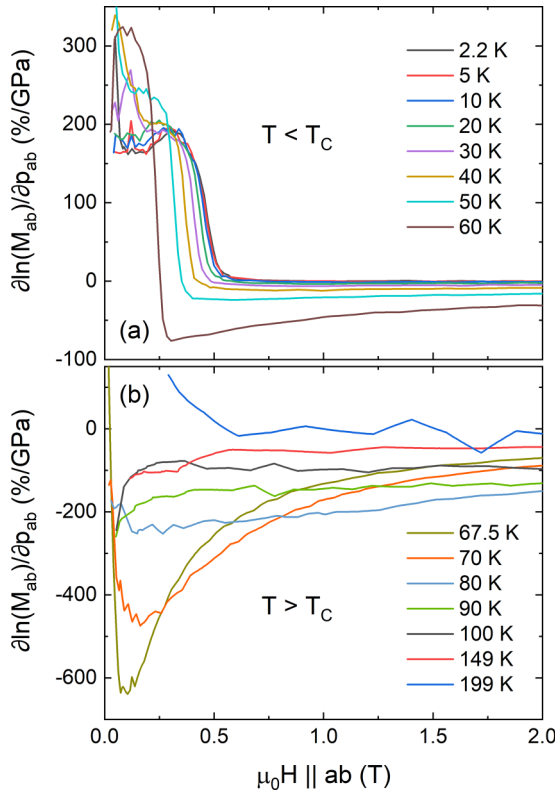


FIG. 4. Pressure dependence of the magnetization for $B, p \parallel ab$ at (a) temperatures $T < T_C$ and (b) $T > T_C$.

In the case at hand $i = j$, i.e., the magnetic field is applied along the direction of the measured length changes.

At $T < T_C$, the in-plane magnetostriction exhibits a positive jump in λ_{ab} at critical fields $B_{cr} \approx 0.5$ T (see Fig. 2) which according to Eq. (9) translates to a negative jump in $\partial \ln(M_{ab})/\partial p_{ab}$. As seen in Fig. 4(a), the uniaxial pressure dependence of the magnetization along ab is large for $B < B_{cr}$, i.e., M_{ab} will be strongly enhanced by uniaxial pressure $p \parallel ab$. In contrast, for $B > B_{cr}$ the pressure effect is close to zero at low temperatures and negative with increasing magnitude as the temperature approaches T_C . Above T_C , $\partial \ln(M_{ab})/\partial p_{ab}$ is strongly negative with a peak at low fields reaching up to $-600\%/GPa$ at 67.5 K [Fig. 4(b)]. As the temperature is increased further above T_C the pressure dependence of the magnetization decreases until it reaches roughly zero at 200 K.

For $B, p \parallel c$, the uniaxial pressure dependence of the magnetization is positive both in the paramagnetic and the ferromagnetic phase (data not shown). Except for the opposite sign, similar behavior is found as observed for $B, p \parallel ab$, i.e., there is a strong pressure effect of several hundred %/GPa just above T_C at low fields whereas the pressure effect amounts to only a few %/GPa at $T < 50$ K and well above T_C .

Neither measurements of the in-plane magnetization of $Cr_2Ge_2Te_6$ under (hydrostatic or uniaxial) pressure nor any measurements at all under uniaxial pressure have been reported in the literature. However, the *hydrostatic* pressure dependence of the magnetization has been measured for $B \parallel c$. Measurements by Sakurai *et al.* at 0.1 T show a pressure dependence of the magnetization of about $-20\%/GPa$ for

pressure up to 3.4 GPa at low temperatures [22]. In contrast, Bhoi *et al.* found only small changes of a few percent up to 1.73 GPa under the same conditions whereas a much stronger pressure dependence also up to about $-20\%/GPa$ has been reported at 4 GPa [38]. Although in vdW materials it may be not straightforward to trace back experimentally observed hydrostatic effects to uniaxial ones, these results suggest that under hydrostatic pressure the resulting strain in the ab plane, which leads to the rotation of the magnetic moments into the ab plane, has a stronger effect than the c axis stress. This result is in line with our observation that $\partial \ln(M_{ab})/\partial p_{ab} \gg \partial \ln(M_c)/\partial p_c$. The here reported uniaxial pressure dependencies are much larger than the hydrostatic ones from the literature which is typical in solids. It indicates mutual canceling of the uniaxial in-plane and out-of-plane pressure dependencies and is in accordance with the cancellation of uniaxial pressure dependencies of the critical temperature in $Cr_2Ge_2Te_6$ [25]. The observed strong uniaxial pressure dependencies around T_C , however, further confirm the relevance and enhancement of critical fluctuations in this regime of the phase diagram.

D. Phase diagrams

From the thermal expansion, magnetostriction, and magnetization measurements, the phase diagrams for $B \parallel c$ and $B \parallel ab$ are derived (Fig. 5). Dashed lines indicate the crossover from the ferromagnetic to the paramagnetic phase. Red arrows show the effects of uniaxial pressure, i.e., the stabilization of the FM low-temperature phase for $B, p \parallel c$, and strong suppression of the LTF phase for $B, p \parallel ab$. Additionally, the sign of the pressure dependence of the magnetization as calculated above is indicated in each of the phases.

Overall, from the anomalies in the strain response functions the “magnetoelastic” phase diagram is constructed as they separate the different phases recently reported in $Cr_2Ge_2Te_6$, i.e., the PM, ferromagnetic LTF and aligned FM phases for $B \parallel ab$ and PM/FM phases for $B \parallel c$ [26,27]. The elastic signatures of the crossover between the PM to the FM phase allows to follow this feature up to high magnetic fields; it displays a left-bending behavior and a moderate field dependence. In addition, anomalous contributions to the thermal expansion and magnetostriction confirm the presence of short-range spin correlations well above T_C . As illustrated by the blue region in Fig. 5(b), full spin alignment driven by magnetic-fields $B \parallel ab$ is associated with a jump in the magnetostriction coefficient, i.e., a kink in the magnetostrictive length changes, which signals a continuous phase transition. Very strong uniaxial pressure dependencies of the magnetization in the LTF phase and of B_{cr} show the relevance of critical behavior in the whole LTF phase. In the vicinity of T_C , different signs of magnetostriction anomalies, the presence of an additional feature in the high-temperature shoulder of the thermal expansion anomaly at small applied magnetic-field $B \parallel ab$ and the suppression of the transition temperature into the LTF phase in magnetic fields clearly supports the scenario of an intermediate ferromagnetic phase, i.e., the fully aligned FM phase, separating the LTF and PM phases at small applied magnetic fields [27]. Our data suggest that the phase boundaries associated with the crossover FM/PM

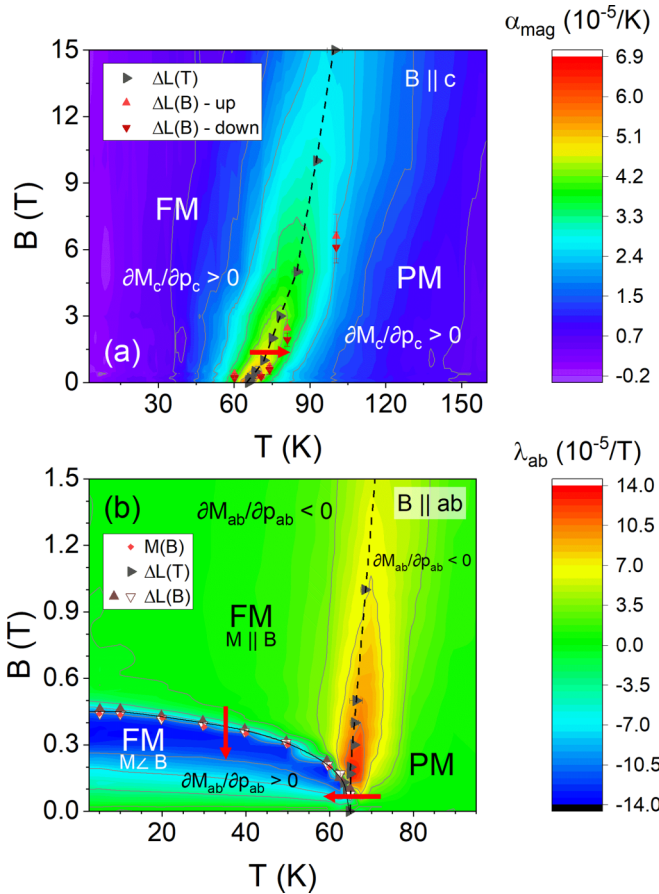


FIG. 5. Phase diagrams for (a) $B \parallel c$ and (b) $B \parallel ab$ with uniaxial pressure dependencies of the magnetization (see the text) as well as of the phase boundaries (red arrows) indicated. Color coding shows the magnetic thermal expansion coefficient (a) as well as the magnetization coefficient (b). Dashed lines indicate the crossover from the FM to the PM phase. The solid black line in (b) indicates the transition from the low temperature and field (LTF) phase to the phase where magnetic moments are fully aligned along the applied field. Red arrows show the effects of uniaxial pressure (see the text).

and the LTF/FM transition, respectively, merge into a triple point which according to Liu *et al.* [31] may be characterized by tricritical behavior. In the vicinity of the triple point, we

observe significantly enhanced spin fluctuations as indicated by strong magnetostrictive response [red and blue regions in Fig. 5(b)] and by very large uniaxial pressure dependencies of the magnetization (see Fig. 4).

V. SUMMARY

Thermal expansion and magnetostriction measurements of the $\text{Cr}_2\text{Ge}_2\text{Te}_6$ single crystals confirm the presence of a low-temperature and low-field (LTF) phase for $B \parallel ab$. Magnetostriction and magnetization data of this phase can be modeled within the Stoner-Wohlfarth model at low temperatures, advocating that as a magnetic-field $B \parallel ab$ is applied the magnetic moments rotate from the c axis towards the applied field until saturation is reached around 0.5 T. The uniaxial pressure dependence of the critical field shows that pressure $p \parallel ab$ strongly suppresses the LTF phase, presumably by lowering the uniaxial anisotropy. This conclusion is further supported by the uniaxial pressure dependence of the magnetization. In the vicinity of T_C , the presence of two distinct thermal expansion anomalies at small fields $B \parallel ab$ of different magnetic-field dependence supports the scenario of an intermediate FM phase separating PM and LTF phases at small fields and is indicative of a triple point. The observed very strong uniaxial pressure dependencies, in particular, in the LTF phase and in the vicinity of the triple point as well as pronounced critical behavior observed in the thermal expansion and the magnetostriction underlines the importance of spin fluctuations in $\text{Cr}_2\text{Ge}_2\text{Te}_6$ and evidences the strong coupling of spin and lattice.

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