



The Interplay of Single Ion Anisotropy and Magnetic 3d-4f Interactions in $V^{III}_2Ln^{III}_2$ Butterfly Complexes

Jan Arneth,^{*,[a]} Xian-Feng Li,^[b] Jonas Braun,^{*,[b, c, d]} Sagar Paul,^[e] Michael Schulze,^[e] Christopher E. Anson,^[b] Wolfgang Wernsdorfer,^[d, e] Annie K. Powell,^{*,[b, c, d]} and Rüdiger Klingeler^{*,[a]}

Within the framework of 3d-4f molecular magnets, the most thoroughly investigated architecture is that of butterfly-shaped coordination clusters as it provides an ideal testbed to study fundamental magnetic interactions. Here, we report the synthesis and characterisation of a series of isostructural $V^{III}_2Ln^{III}_2$ butterfly complexes, where $Ln=Y$ (1_Y), Tb (2_{Tb}), Dy (3_{Dy}), Ho (4_{Ho}), Er (5_{Er}), Tm (6_{Tm}), Yb (7_{Yb}), which extends the previous study on isostructural butterflies with Cr^{III} , Mn^{III} and Fe^{III} . In zero external field, compounds 2_{Tb} , 3_{Dy} and 4_{Ho} show clear maxima in the out-of-phase component of the ac susceptibility whereas small magnetic fields are needed to suppress quantum

tunnelling in 6_{Tm} . Combined high-field electron paramagnetic resonance spectroscopy and magnetisation measurements unambiguously reveal an easy-plane anisotropy of the V^{III} ion and antiferromagnetic Ising-like 3d-4f exchange couplings. The strength of J_{3d-4f} is shown to decrease upon variation of the 4f ion from Tb to Ho, while increasing antiferromagnetic interaction can be observed from Ho to Tm. The exact inverse chemical trend is found for the relative angle between the 3d and 4f main anisotropy axes, which highlights the important role of the lanthanide 4f electron distribution anisotropy for 3d-4f exchange.

Introduction

One approach to enhance the performance of Single Molecule Magnets (SMMs) such as coordination clusters which exhibit magnetic hysteresis solely due to magnetisation blocking on the molecular level,^[1] is the idea of designing heterometallic 3d-4f complexes. This emerged as a promising pathway to higher blocking temperatures (T_B).^[2-4] In these compounds the magnetic relaxation of the strongly anisotropic lanthanide moments, which is typically affected by quantum tunnelling and spin-phonon based relaxation pathways, can be slowed down via favourable magnetic interactions with the transition metal ions.^[5-7] Although the combination of high effective energy barriers (U_{eff}) in 4f-based systems and slow magnetic relaxation of exchange coupled 3d ion systems appears to be a logical strategy, only a few examples in which the presence of paramagnetic 3d moments significantly improves the SMM properties have been reported.^[8-10] Clearly the magnetic exchange coupling J_{3d-4f} is one of the most important parameters influencing the performance of 3d-4f-based magnets. Hence knowledge on the underlying microscopic mechanisms is crucial for further advancements. The electronic structure of 4f ions is a further complication in the determination of the size and sign of J_{3d-4f} and predictions are mostly based on quantum chemical calculations.^[8,10-12] On the experimental side, high-frequency/high-field electron paramagnetic resonance (HF-EPR) spectroscopy has established itself as the main technique to investigate the magnetic 3d-4f coupling in heterometallic complexes as it allows to quantify the coupling parameter J_{3d-4f} precisely.^[13-16] Other spectroscopic methods, such as inelastic neutron scattering,^[17] magnetic circular dichroism^[18] and far-infrared transmission spectroscopy,^[19] give further insights into the microscopic parameters governing the magnetic behaviour

[a] J. Arneth,^{*} Prof. Dr. R. Klingeler
Kirchhoff-Institute for Physics
Heidelberg University
INF 227, D-69120 Heidelberg, Germany
E-mail: jan.arneth@kip.uni-heidelberg.de
ruediger.klingeler@kip.uni-heidelberg.de

[b] Dr. X.-F. Li,^{*} Dr. J. Braun, Dr. C. E. Anson, Prof. Dr. A. K. Powell
Institute of Inorganic Chemistry
Karlsruhe Institute of Technology (KIT)
Kaiserstr. 12, D-76131 Karlsruhe, Germany
E-mail: jonas.braun2@kit.edu
annie.powell@kit.edu

[c] Dr. J. Braun, Prof. Dr. A. K. Powell
Institute of Nanotechnology
Karlsruhe Institute of Technology (KIT)
Kaiserstr. 12, D-76131 Karlsruhe, Germany
E-mail: jonas.braun2@kit.edu
annie.powell@kit.edu

[d] Dr. J. Braun, Prof. Dr. W. Wernsdorfer, Prof. Dr. A. K. Powell
Institute for Quantum Materials and Technology
Karlsruhe Institute of Technology (KIT)
Kaiserstr. 12, D-76131 Karlsruhe, Germany
E-mail: jonas.braun2@kit.edu
annie.powell@kit.edu

[e] Dr. S. Paul, Dr. M. Schulze, Prof. Dr. W. Wernsdorfer
Physikalisches Institut
Karlsruhe Institute of Technology (KIT)
Engesserstr. 15, D-76131 Karlsruhe, Germany

[[†]] These authors have contributed equally.

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of 3d–4f magnets. Among the vast variety of available 3d–4f architectures, the butterfly-motif provides a useful test-bed system and many systematic studies have been performed regarding both their static and dynamic magnetic properties.^[20–24] This motif can either have the 3d metal ions in the central “body” position with the 4f ions at the “wingtips” (Type I) or the inverse situation for Type II. The investigation of V^{III}-based butterfly systems has not been reported so far because of difficulties in obtaining stable V–Ln clusters with this desired oxidation state.

Here, we report the successful synthesis using Schlenk line techniques of a series of isostructural V^{III}₂Ln^{III}₂ Type II butterfly complexes, where Ln = Y (1_Y), Tb (2_{Tb}), Dy (3_{Dy}), Ho (4_{Ho}), Er (5_{Er}), Tm (6_{Tm}), Yb (7_{Yb}). Using combined high-field EPR and magnetisation studies we quantitatively determined the microscopic spin Hamiltonian parameters, such as 3d–4f coupling and relative angles between the main anisotropy axes. The dependence of J_{3d-4f} on the 4f ion is discussed and compared to the experimental findings on other 3d–4f complexes. We show that two distinct chemical trends can be observed for lanthanide ions with predominantly oblate and prolate anisotropy shape and conclude that the 4f electron distribution plays a crucial role in the dominant 3d–4f charge transfer process. The importance of 4f charge density anisotropy becomes further evident from the magnetic relaxation behaviour as studied by ac susceptibility and low temperature μ SQUID measurements.

Experimental Detail and Methods

Single crystal crystallographic data were measured on Stoe StadiVari diffractometers equipped with Mo- or Cu-microfocus sources or a MetalJet2 liquid Ga source. Structures were solved using SHELXT^[25] and refined using SHELX-2019^[26] within the Olex2 platform.^[27] The lattice methanols and the methyls of some ligand *t*-butyl groups were disordered and refined with sets of partial occupancy atoms. These were assigned anisotropic thermal parameters except for minor components, and refined with similarity restraints applied to bond lengths and rigid-bond restraints applied to the temperature factors. Crystallographic data are given in Table S1. Full crystallographic data and details of the structural determinations for the structures in this paper have been deposited with the Cambridge Crystallographic Data Centre as supplementary publication nos. CCDC 2406392–2406398. Copies of the data can be obtained, free of charge, from <https://www.ccdc.cam.ac.uk/structures/>.

The direct current (dc) and alternating current (ac) magnetisation was studied in the temperature range $T = 1.8$ –300 K by means of a Magnetic Properties Measurement System (MPMS3, Quantum Design) and a Physical Properties Measurement System (PPMS, Quantum Design). For all measurements powder samples were pelletised in polycarbonate capsules to avoid reorientation in external magnetic fields. The experimental data were corrected for the diamagnetic contribution of the sample holder and of the ligands calculated by means of Pascal's

constants.^[28] Simulations of the magnetic data were performed using the PHI software package.^[29]

$M(B)$ measurements over a temperature range of 0.03–5.0 K, on single crystals of compound 3_{Dy}, were conducted using μ SQUIDS. A single crystal of about 100 μ m (long arm) length was placed near an array of μ SQUIDS on a chip, within few μ m gap between the crystal edge and the μ SQUID loop to ensure optimal coupling. The crystal was thermalized with ApiezonTM grease and cooled to a base temperature of 30 mK in a dilution refrigerator. A 3D vector magnet allows for the application of the magnetic field in different directions within the SQUID plane with an angular accuracy better than 0.1°. Low-temperature $M(B)$ measurements were carried out at varying field sweep rates (0.001 to 0.128 T s^{−1}) with a time resolution of ~ 1 ms. The easy axis of the crystal was identified using the ‘transverse field method’.^[30–31]

High-frequency/high-field electron paramagnetic (HF-EPR) resonance studies were performed using a phase sensitive millimetre vector network analyser (MVNA) by ABmm as a simultaneous microwave source and detector.^[32] Temperature control from 2 K to 300 K was ensured by placing the sample space in the Variable Temperature Insert (VTI) of an Oxford magnet system equipped with a 16 T superconducting coil.^[33] Polycrystalline powder samples were prepared in a brass ring sealed with kapton tape either as loose powder or fixed by mixing with eicosane. The former setup allows alignment of the crystallites with the external magnetic field, hence providing simplified pseudo-single-crystal spectra.^[16,34–37] Alignment is ensured by sweeping up the magnetic field to 16 T prior to each measurement and confirmed by observation of corresponding orientation jumps in the transmitted microwave intensity signal. Spectral simulations of the HF-EPR data were performed using the EasySpin software package.^[38]

Experimental Results

Structural Analysis

Compounds 1_Y to 7_{Yb} with the formula [V₂Ln₂(μ_2 -OMe)₂(Budea)₂(piv)₆] \cdot 2CH₃OH (Ln^{III} = Y^{III}, Tb^{III}, Dy^{III}, Ho^{III}, Er^{III}, Tm^{III} and Yb^{III}) all crystallise isomorphously and thus isostructurally in the triclinic space group P $\bar{1}$ with $Z = 1$ (Table S1); the molecular structure of 3_{Dy} is shown in Figure 1. The centrosymmetric V₂Dy₂ core corresponds to a Type II butterfly structure^[24] with Dy(1) and Dy(1') occupying the body position and V(1) and V(1') are located at the wingtips. The two VDy₂ triangles are each bridged by a triply-bridging methoxide ligand, one above and the other below the V₂Dy₂ plane. The vanadiums are both chelated by a doubly-deprotonated (Budea)^{2−} ligand with the two oxygens each bridging to one of the Dy^{III} ions. The coordination is completed by four *syn,syn*-pivalates bridging the V–Dy edges of the core with two further pivalates each chelating a Dy^{III} ion.

The clusters are isostructural to Cr₂Dy₂ butterflies reported by Murray *et al.*^[8] Selected bond lengths and angles are compared in Table S2, where the expected trends in the Ln–O

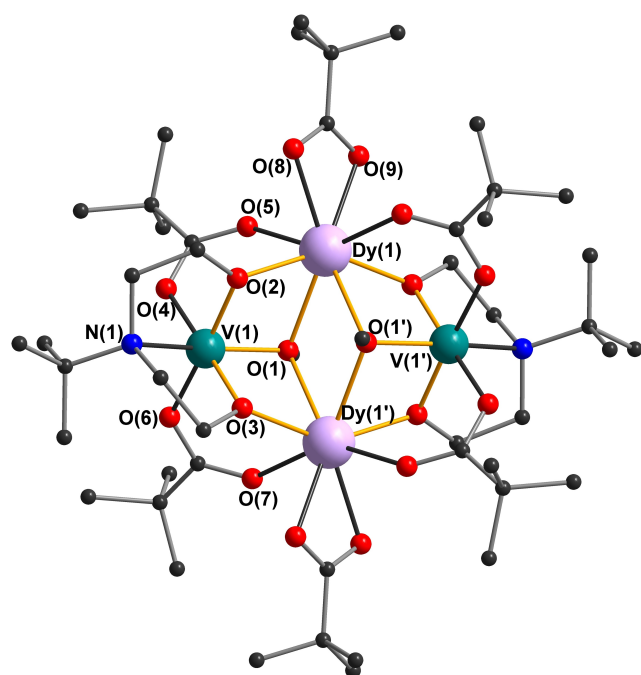


Figure 1. Molecular structure of the V_2Dy_2 cluster in $3Dy$. Hydrogen atoms and minor disorder are omitted for clarity. Primed atoms at 1-x, 1-y, 1-z.

distances given the change in size of the Ln^{III} ions is clear. The V^{III} oxidation states in the complexes were confirmed by Bond Valence Sum Analysis,^[39] with the calculated valences for $V(1)$ being in the range 2.94–2.97 (see Table S3).

Direct Current (dc) Magnetisation

Variable-temperature static magnetic susceptibility measurements performed at $B=0.1$ T on polycrystalline powders of 1_V – 7_{Yb} result in the χT -curves shown in Figure 2. For all the measured compounds the room temperature (300 K) χT -values which are summarized in Table 1 are slightly smaller than the expected free ion values for an uncoupled system of two V^{III} ions with $g=2$ and two Ln^{III} ions. Such an observation is not surprising, since the LS-coupling in V^{III} typically leads to g -factors smaller than 2.^[40–45] In the case of 1_V , the magnetic data imply an average $g_{ave}=1.93$, which is in excellent agreement with the results of our HF-EPR measurements, as discussed later. However, the influence of spin-orbit coupling in V^{III} can only partly explain the lower experimentally determined room temperature χT of 2_{Tb} and 3_{Dy} . In these compounds, the too low value of χT might be the result of solvent molecules remaining in the crystal matrix or from the strongly anisotropic nature of the Ln moments such that even at 300 K not all m_J sub-levels are equally populated.^[46]

The χT -product of 1_V is almost constant at temperatures above 10 K and rapidly decreases upon cooling below this value. As will be revealed by the HF-EPR studies, this drop is mainly attributed to the depopulation of the m_J sub-levels, in addition to a weak antiferromagnetic interaction between the

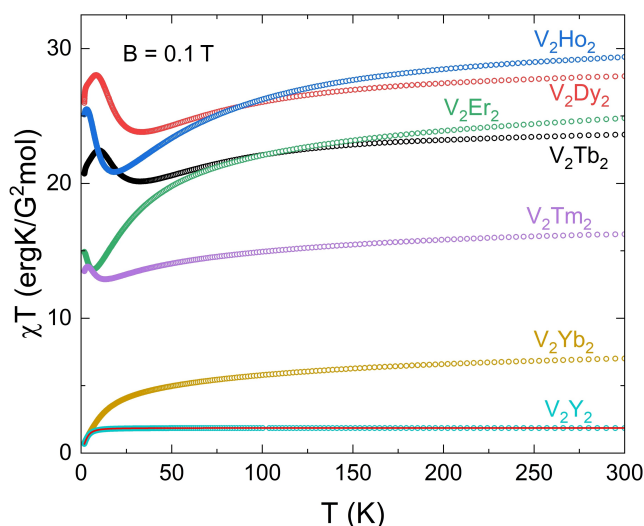


Figure 2. Temperature dependence of the χT -product of 1_V – 7_{Yb} in an external magnetic field of $B=0.1$ T. The red solid line depicts a simulation of 1_V using the parameters obtained from HF-EPR (see text).

Table 1. Experimentally determined and theoretical free ion χT -values (in $ergK/G^2mol$) of 1_V – 7_{Yb} .

	g_{Ln}	J_{Ln}	χT_{exp} (1.8 K)	χT_{exp} (300 K)	$\chi T_{free ion}$
1_V			0.68	1.86	2.0
2_{Tb}	3/2	6	20.7	23.6	25.6
3_{Dy}	4/3	15/2	26.0	28.0	30.3
4_{Ho}	5/4	8	25.1	29.4	30.1
5_{Er}	6/5	15/2	14.9	24.8	25.0
6_{Tm}	7/6	6	13.5	16.2	16.3
7_{Yb}	8/7	7/2	0.76	7.0	7.14

V^{III} ions. A similar temperature dependence of χT is found in 7_{Yb} , albeit a weak decrease already begins close to room temperature and the steep drop occurs already at around 25 K. Both signatures presumably arise from the depopulation of the strongly split m_J sub-levels in the anisotropic Yb^{III} ions. A priori, antiferromagnetic intramolecular Yb – Yb coupling or intermolecular interaction might also yield such behaviour. However, for both we estimate an upper limit of dipolar interaction energy of less than 0.3 K which is way too small to account for the observed temperature dependence.

In contrast, the $\chi T(T)$ -profiles for 2_{Tb} – 6_{Tm} all exhibit a maximum at temperatures in the range 2–10 K, superposed upon the general decrease on cooling. Such behaviour could be attributed to intramolecular ferromagnetic interactions between the V^{III} and the Ln^{III} moments. However, the HF-EPR measurements reveal that it can actually be attributed to a parallel arrangement of the Ln^{III} spins on opposite sides of the butterfly body mediated by the antiferromagnetic V^{III} – Ln^{III} interactions. The position of the maximum in χT gives information about the J_{V-Ln} coupling strength, as discussed below.

High-Field EPR Studies

HF-EPR measurements on freshly ground powder samples from the same batch were carried out to determine the microscopic magnetic parameters. Figure 3 shows representative HF-EPR spectra of 1_V at variable temperature and fixed microwave frequency of $f=395.4$ GHz (a) as well as at variable frequency and fixed temperature of $T=2$ K (b). The spectra exhibit clear resonances in the accessible field and frequency range. In total two asymmetric resonance features can be identified as indicated by the vertical dashed lines in Figure 3(a). Upon increasing temperature both features become less pronounced, indicating that the initial state of the corresponding transition is (at least close to) the magnetic ground state. Moreover, the corresponding frequency dependence suggests that the low-field feature is associated with a symmetry-forbidden transition yielding $\Delta m=2$, while the high-field feature is related to an allowed transition with $\Delta m=1$. The allowed high-field feature

exhibits its main spectral weight on its high-field side, thereby directly implying an easy-axis-type g -factor anisotropy.^[47–48]

In order to quantify the qualitatively discussed parameters we performed spectral simulations using an interacting dimer Hamiltonian in the form of

$$\hat{\mathcal{H}}_{1_V} = \sum_{i=1}^2 \left(\mu_B \mathbf{B} \hat{g} \hat{S}_i + D \hat{S}_{i,z}^2 + E (\hat{S}_{i,x}^2 - \hat{S}_{i,y}^2) \right) - J \hat{S}_1 \hat{S}_2 \quad (1)$$

where $S_1=S_2=1$, $\vec{g}=(g_{\perp}, g_{\perp}, g_{\parallel})$ denotes an axial g -tensor, D and E are the axial and rhombic anisotropy constants and J represents an isotropic Heisenberg interaction between the V^{III} moments. Note here that both the g -tensor and the crystal field parameters are fixed to be equal for both V^{III} sites, since their local ligand coordination, and hence the surrounding electrostatic potentials, are equivalent under point inversion. The best simulations of the experimental data were achieved by the parameters $g_{\perp}=1.91(2)$, $g_{\parallel}=1.98(2)$, $D=+9.3(4)$ K (6.45 cm^{-1}), $|E|=0.25(5)$ K (0.17 cm^{-1}), $J=-0.18(6)$ K (0.12 cm^{-1}) and are shown as red solid lines in Figure 3. The spectral simulations confirm the easy-axis-character of the g -tensor, while simultaneously revealing an easy-plane single-ion anisotropy with a small rhombic distortion. The magnitude of the D and E parameters is not uncommon for V^{III} compounds reported in the literature.^[40–45] In addition, this HF-EPR analysis reveals a weak antiferromagnetic coupling between the V^{III} moments which likely arises from a superexchange interaction *via* the diffuse empty orbitals of the diamagnetic Y^{III} ion.^[49] Although D is the dominant term it is necessary to include the contribution of the parameters E and J in order to simulate the experimental data well. The parameter set extracted from HF-EPR also fully reproduces the observed magnetic data of 1_V , as can be seen by the simulation depicted in Figure 2.

We additionally performed HF-EPR spectroscopy on loose powder samples of $2_{\text{Tb}}7_{\text{Vb}}$, for which 2_{Tb} , 5_{Er} and 6_{Tm} show well-resolved resonance features in the accessible frequency- and field-range. The obtained resonance positions for $T=2$ K are summarised in the magnetic field vs. resonance frequency diagrams in Figure 4, where distinct resonance features are marked by different symbols. Compared to the spectra of 1_V , the resonance features in all three compounds are significantly less pronounced and exhibit strongly frequency-dependent line shapes. Upon following the main spectral weight clear resonance branches can be identified.

In 2_{Tb} a single EPR-transition can be observed at frequencies ranging from 280 to 500 GHz. Up to at least 8 T its resonance frequency decreases linearly with magnetic field, corresponding to an effective g -factor of $g_{\text{eff}} \approx 2$, close to the g -value found for the V^{III} ions in 1_V . This suggests that the transition detected here is likely to be associated with an excitation involving the V^{III} ions in 2_{Tb} . However, the data also indicate a significant magnetic coupling to the Tb^{III} moments, since the extrapolated zero-field excitation gap $\Delta \approx 515$ GHz of the resonance branch cannot be explained by the single-ion anisotropy of the V^{III} ions alone. The low temperature HF-EPR spectra of 5_{Er} and 6_{Tm} exhibit up to two (5_{Er}) and four (6_{Tm}) resonance features yielding

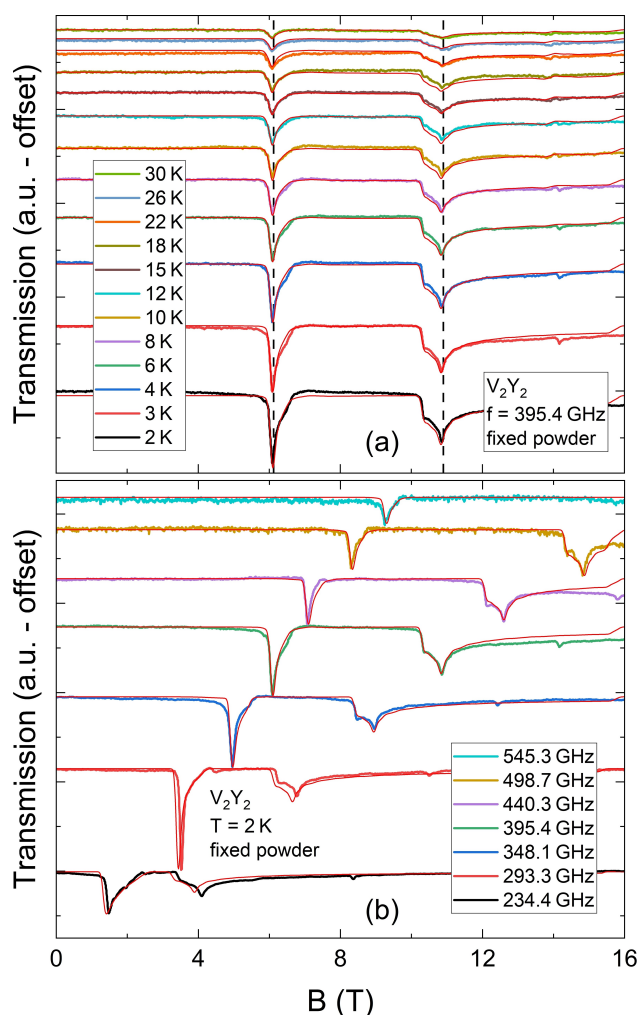


Figure 3. HF-EPR spectra of a fixed powder sample of 1_V at different temperatures and fixed frequency $f=395.4$ GHz (a) and at different frequencies and constant temperature $T=2$ K (b). Vertical dashed lines indicate the distinct resonance features. Red solid lines depict simulations using equation 1 with the parameters described in the main text.

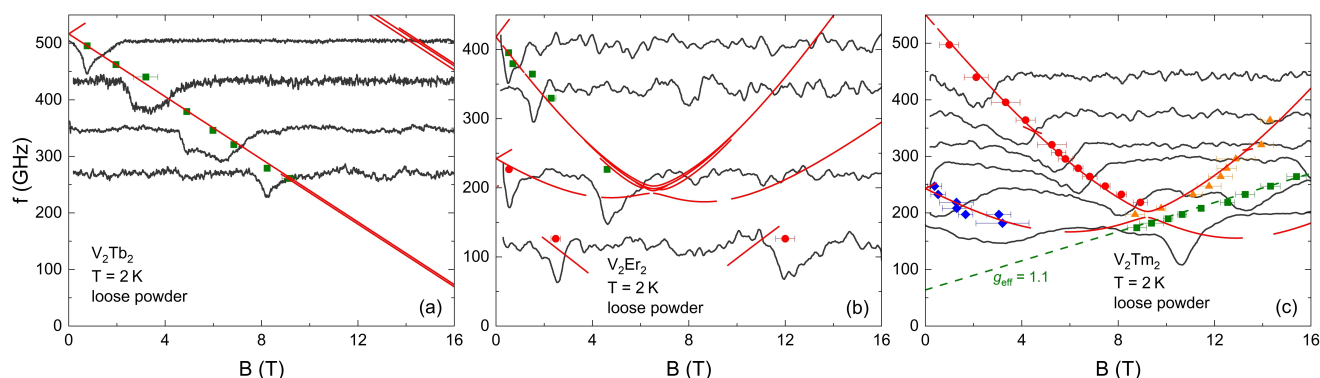


Figure 4. Magnetic field vs. resonance frequency plot of 2_{Tb} (a), 5_{Er} (b) and 6_{Tm} (c) obtained from HF-EPR measurements at $T = 2$ K. Red solid lines depict simulations using equation 2 with the parameters shown in Table 2. Selected HF-EPR spectra are visible in the background and are shifted along the ordinate to align with the corresponding measurement frequency. The green dashed line in (c) marks the resonance branch which is associated with a transition between different m_j levels of the Tm^{III} ions as described in the text.

strongly curved branches with initial slopes corresponding to $g_{\text{eff}} \neq 2$. Interestingly, both complexes share a common low-energy zero-field excitation gap at $\Delta \simeq 240$ GHz, whereas the high-energy gaps differ significantly from each other. The temperature dependence of the HF-EPR spectra indicates that all observed transitions arise from either the ground or low-lying excited states (see Supplementary Material).

As a result of their large total angular momentum and complex anisotropy shape of Ln^{III} ions, simulations using full spin Hamiltonians for these ions in low-symmetry ligand fields generally exceed the manageable Hilbert and parameter space. However, in suitable coordination environments an axial charge distribution can lead to well-isolated ground state doublets, which allow for a treatment of Ln^{III} moments as Ising-spins at sufficiently low temperatures.^[15–16,50–51] In the frame of this model the Hamiltonian of $2_{\text{Tb}}-7_{\text{Vb}}$ reads

$$\hat{H}_{V_2\text{Ln}_2} = \vec{R}_x^T(\theta) \hat{H}_{1_V} \vec{R}_x(\theta) + \sum_{i=3}^4 \left(\mu_B 2\alpha \vec{B} \vec{g} \vec{S}_i - 2\alpha J_{V-\text{Ln}} (\hat{S}_{1,z} \hat{S}_{i,z} + \hat{S}_{2,z} \hat{S}_{i,z}) \right) \quad (2)$$

where $S_3 = S_4 = 1/2$, $\vec{g} = (0, 0, g_{\text{Ln}})$, α is a scaling factor to account for renormalization in the pseudospin transformation, and $\vec{R}_x(\theta)$ denotes a rotation of the local 1_V -subsystem parameters (\vec{g} and \vec{D}) by θ around the global x-axis.

Since the local environment around the V^{III} ions is essentially identical in all compounds, it is justifiable to fix the parameters obtained from the HF-EPR studies on 1_V in the simulations of $2_{\text{Tb}}-7_{\text{Vb}}$. In addition, it is assumed that the strongly Ising-like character of the Ln^{III} moments means that in a loose powder sample the crystallites can be oriented by the applied field such that the easy-anisotropy axes coincide with that of the lanthanide moments (here chosen parallel to the z-axis) and, hence, leads to an alignment of the molecules along the z-axis. The best simulations of the resonance branches are achieved by using $2\alpha J_{V-\text{Tb}} = -34.1$ K and $\theta = 0^\circ$ for 2_{Tb} , $2\alpha J_{V-\text{Er}} = -8.6$ K and

$\theta = 60^\circ$ for 5_{Er} , and $2\alpha J_{V-\text{Tm}} = -12.2$ K and $\theta = 45^\circ$ for 6_{Tm} , respectively. A more elaborate overview of the obtained spin Hamiltonian parameters is given in Table 2. It should be noted that neither the inclusion of magnetic dipolar coupling between the Ln^{III} moments nor an additional rotation of the V^{III} ion subsystem around its local y-axis significantly improved the simulations.

As can be seen in Figure 4, the experimentally observed resonance branches are well described by the spin Hamiltonian in equation 2. Only the low-energy branch at high magnetic fields in 6_{Tm} (green squares) cannot be reproduced by our model. The slope of this branch corresponds to an effective g -factor of $g_{\text{eff}} \simeq 1.1$, which is close to the Landé-factor of Tm^{III} . Therefore, it seems reasonable that this resonance feature is associated with a transition between different m_j levels of the Tm^{III} ions that are effectively neglected in the Ising-approximation used here. Since α and $J_{V-\text{Ln}}$ are linearly correlated (see equation 2), this analysis of the HF-EPR data does not allow for a determination of the coupling constant $J_{V-\text{Ln}}$ but only of the magnetic coupling energy $2\alpha J_{V-\text{Ln}}$. In order to decouple the two parameters, we simulated the low-temperature magnetic data using the parameters obtained from the HF-EPR studies.

Figure 5 shows the isothermal magnetisation of 2_{Tb} at $T = 1.8$ K and the corresponding simulations with $\alpha = 6$ and, hence, $J_{V-\text{Tb}} = -2.84$ K. The chosen model is in line with the experimental data up to at least 6 T and reproduces the key characteristics of the $M(B)$ -profile *i.e.* a steep increase at small magnetic fields and an inflection point, indicating a change in the magnetic ground state, at $B \simeq 5$ T. The discrepancy at higher magnetic

Table 2. Spin Hamiltonian parameters obtained from combined analysis of HF-EPR and magnetisation data using equation 2.

	α	$J_{V-\text{Ln}}$ (K)	\vec{g}_L	θ ($^\circ$)
2_{Tb}	6	−2.84(4)	0.32	0(4)
4_{Ho}	7	−0.55(11)	0.1	85(16)
5_{Er}	6	−0.72(7)	0.32	60(8)
6_{Tm}	6	−1.02(6)	0.15	45(5)

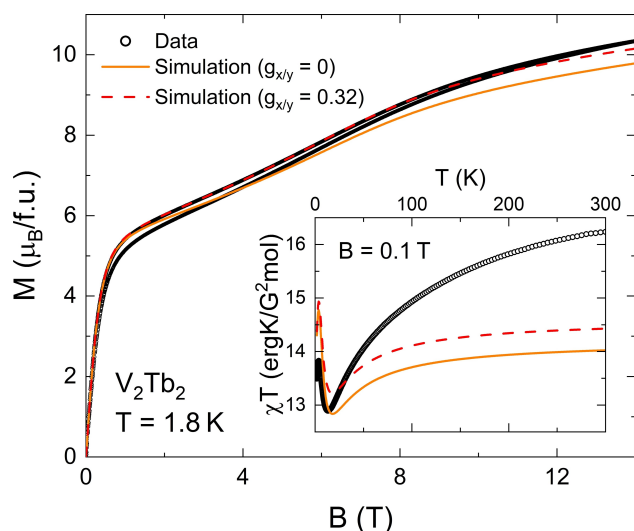


Figure 5. Isothermal magnetisation loop (0 T → 7 T → 0 T) of V_2Tb_2 at $T = 1.8$ K and simulations of the down-sweep using the strict Ising model (orange solid line) in equation 2 and an XXZ model (red dashed line) with the parameters shown in Table 2 (see the text). The inset depicts the measured and simulated $\chi T(T)$ at $B = 0.1$ T.

fields is likely associated with finite contributions from the transverse components of the Ln^{III} moments, *i.e.* their deviation from perfect Ising-spins in the fixed powder measurements. This can be modelled by introducing transversal components of the effective lanthanide g -tensor, such that $\vec{g} = (\tilde{g}_\perp, \tilde{g}_\perp, g_\parallel)$. As seen in Figure 5, the experimental data are well simulated using this extended model with $\tilde{g}_\perp = 0.32$ over the entire magnetic field range.

Whereas the low-temperature χT -product of V_2Tb_2 (inset of Figure 5) can be fitted by the simulations using the same parameters as used for the $M(B)$ simulation, the Hamiltonian in equation 2 is not able to describe the χT -product at elevated temperatures. This is not unsurprising since the population of higher-lying m_J sub-levels means that the Ln^{III} moments no longer equate to Ising-spins. However, the qualitative behaviour of $\chi T(T)$, especially the position of the characteristic maximum at low and minimum at intermediate temperatures, is in line with this model.

An analogous analysis of combined HF-EPR and magnetisation data on V_2Er_2 and V_2Tm_2 (see Figure S2) results in the spin Hamiltonian parameters given in Table 2. In addition, despite showing no detectable HF-EPR signal, simulations of the magnetisation data also allow the determination of the microscopic parameters in V_2Ho_2 , albeit with reduced accuracy.

Alternating Current (ac) Susceptibility

In order to probe their magnetic relaxation dynamics, ac susceptibility measurements with an oscillating field of $H_{\text{ac}} = 3$ Oe were performed on all investigated complexes. Compounds V_2Tb_2 , V_2Dy_2 and V_2Ho_2 exhibit detectable out-of-phase (χ'') signals without applied dc field. Field-induced slow relaxation of the

magnetisation with an applied dc field of 0.1 T could be observed for V_2Tm_2 , suggesting that quantum tunnelling is the dominant relaxation mechanism when no external field is present. For all other compounds no χ'' signals could be observed in the accessible field and frequency range. The measured frequency dependence of the out-of-phase susceptibilities are shown in Figure 6. Apart from the measurement of V_2Dy_2 two distinct relaxation processes can clearly be discerned for all other compounds.

To corroborate this assignment, the data were fitted using a generalized Double-Debye-model, *i.e.* the sum of two Debye-relaxation functions

$$\chi_{\text{ac}}(f) = \sum_{n=1}^2 \left(\chi_{\text{S},n} + \frac{\chi_{\text{T},n} - \chi_{\text{S},n}}{1 + (if\tau_n)^{1-\alpha_n}} \right) \quad (3)$$

where χ_{S} and χ_{T} are the adiabatic and isothermal susceptibilities, τ denotes the relaxation time, and $0 < \alpha < 1$ qualitatively represents the distribution of relaxation times in the molecule. As can be seen from Figure 6 the chosen model describes the experimental results well. We note that one Debye function is not sufficient to reproduce the ac susceptibility of V_2Dy_2 . This implies the presence of at least two relaxation processes with similar relaxation times. The nature of the underlying microscopic mechanisms leading to slow magnetic relaxation can be elucidated by the temperature dependence of τ . Qualitatively, the fast relaxation processes in V_2Tb_2 , V_2Ho_2 and V_2Tm_2 shift to higher frequencies with increasing temperature, while the broad maximum in $\chi''(f)$ of V_2Dy_2 is almost temperature independent, indicative of dominant quantum tunnelling.

The relevance of under-barrier relaxation in V_2Dy_2 can also be seen from the μSQUID measurements depicted in Figure 7. At sub-Kelvin temperatures the single crystal $M(B)$ loops show a butterfly-shaped hysteresis which indicates fast zero field QTM and a slow thermal relaxation process in finite external

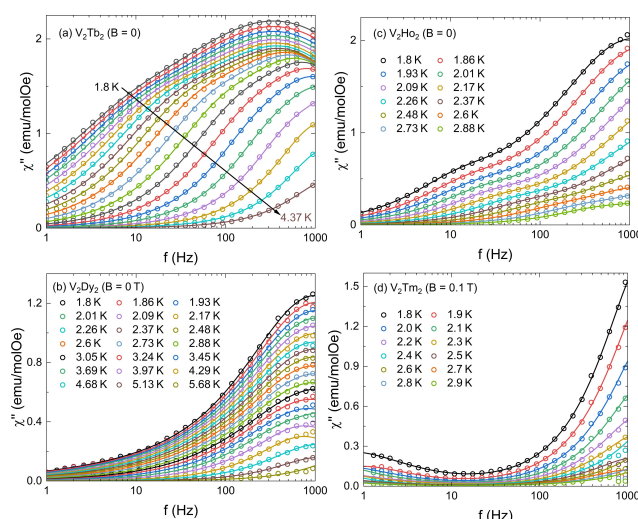


Figure 6. Out-of-phase ac susceptibility of V_2Tb_2 (a), V_2Dy_2 (b), V_2Ho_2 (c) and V_2Tm_2 (d) at selected temperatures. Solid lines depict fits of equation 3 to the data as described in the text. Corresponding in-phase ac susceptibility data and Cole-Cole plots are shown in the Supplementary Material.

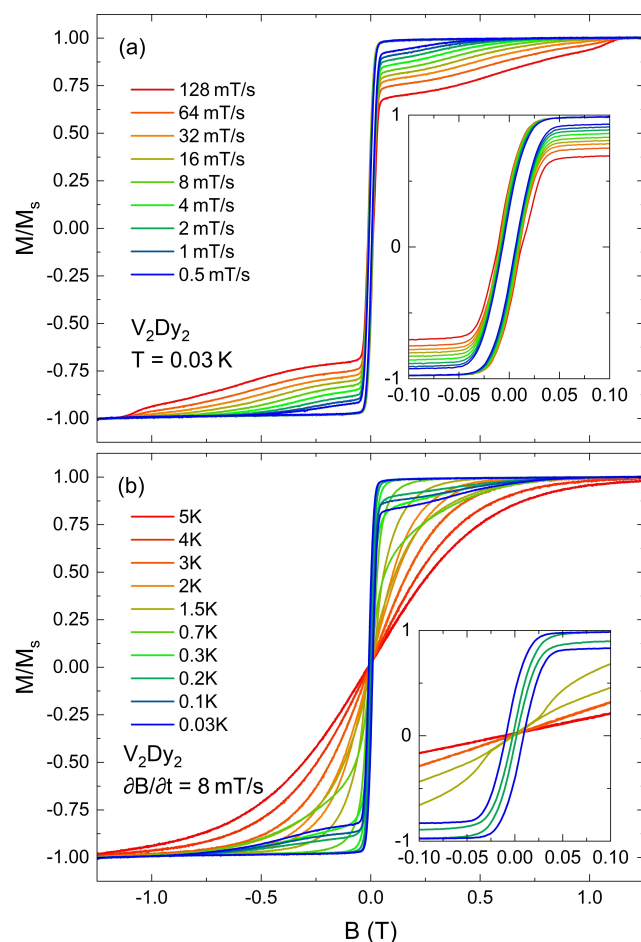


Figure 7. $M(B)$ measurements on a single crystal of $\mathbf{3}_{\text{Dy}}$, at fixed bath temperature $T = 0.03$ K and different sweep rates of magnetic field (a), and fixed sweep rate $\partial B/\partial t = 8$ mT/s and different bath temperatures (b). The insets show the enlarged version of corresponding figures considering only a few curves for clarity.

magnetic fields. In agreement with our ac susceptibility data, even the fastest field scan rate of $\partial B/\partial t = 128$ mT/s yields only a small remanent magnetisation (Figure 7(a)). However, the magnetisation loops remain open (Figure 7(a) inset) with a small coercive field for all sweep rates. SMM behaviour in $\mathbf{3}_{\text{Dy}}$ at such low temperatures probably arises as a result of small effective internal fields originating from intramolecular magnetic interactions. Upon heating from the lowest temperatures (Figure 7(b)), the hysteresis region initially shrinks as it is expected when the contribution of temperature dependent relaxation processes increases. Interestingly, the inverse effect, i.e., widening of the hysteresis on increasing the temperature, can be observed in the temperature range $0.2 \text{ K} < T < 0.9 \text{ K}$. This is not a rare observation in low temperature $M(B)$ measurements when QTM and thermal relaxation times compete at an avoided level crossing region.^[52–53] While the overall shape of the $M(B)$ loops at 30 mK is consistent with a predominantly ferromagnetic character, a magnetisation plateau appears near zero field as the bath temperature is increased (cf. $M(B, T = 1.5 \text{ K})$ in the inset of Figure 7(b)). Using the exchange field approach,^[53] an antiferromagnetic Dy–Dy coupling of $J_{\text{Dy–Dy}} =$

$-6(1) \text{ mK}$ can be read off the data (for a more detailed analysis and discussion see the SI).

A more quantitative investigation of the magnetic relaxation can be provided for both relaxation processes in $\mathbf{2}_{\text{Tb}}$ and for the slow process in $\mathbf{4}_{\text{Ho}}$ from a plot of τ vs T (cf. Figure 8) using the relaxation times extracted from the Debye fits. The linearity of $\ln(\tau)$ vs $1/T$ at high temperatures implies the presence of Orbach and/or slow Raman relaxation. In contrast, faster relaxation mechanisms, such as quantum tunnelling or fast Raman relaxation, dominate at low temperatures. We fit our data over the whole temperature range using

$$\tau^{-1} = \tau_{\text{QTM}}^{-1} + \tau_{\text{R}}^{-1} \frac{e^{-W/k_{\text{B}}T}}{(1 + e^{-W/k_{\text{B}}T})^2} + \tau_{\text{O}}^{-1} e^{-U_{\text{O}}/k_{\text{B}}T} \quad (4)$$

Here, the first term denotes the temperature independent quantum tunnelling of magnetisation, the second term describes Raman relaxation involving only energetically low-lying optical phonons as suggested independently by Gu and Wu and by Lunghi *et al.*^[54–56] and the third term represents Orbach relaxation. The resulting fit parameters, summarised in Table 3, indicate that at low temperatures quantum tunnelling is the main contribution to the fast relaxation process in $\mathbf{2}_{\text{Tb}}$, whereas the slower process is dominated by slow Raman relaxation. This finding is further corroborated by ac susceptibility measurements under a 0.1 T dc field showing that only one relaxation process is visible, which is mainly governed by Orbach and Raman relaxation (see Supplementary Material). At

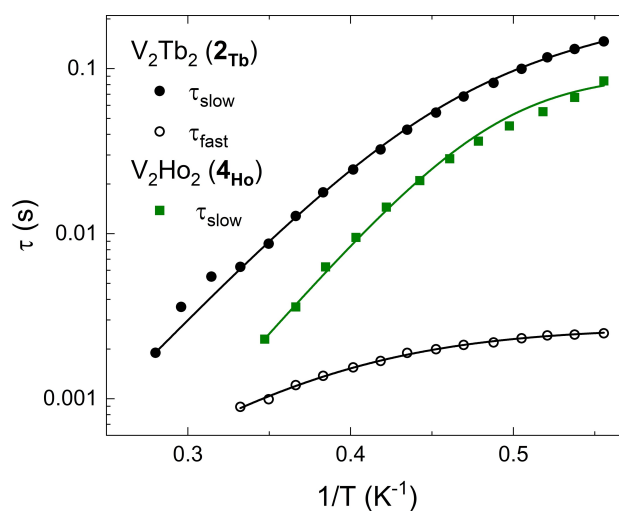


Figure 8. Temperature dependence of the relaxation times in $\mathbf{2}_{\text{Tb}}$ and $\mathbf{4}_{\text{Ho}}$ in the absence of an external static magnetic field. Solid lines depict fits to the data using equation 4 as described in the text.

Table 3. Parameters obtained by the fit of the temperature dependence of the relaxation times in $\mathbf{2}_{\text{Tb}}$ and $\mathbf{4}_{\text{Ho}}$ using equation 4.

	τ_{QTM} (s)	τ_{R} (s)	W (K)	τ_{O} (s)	U_{O} (K)
$\mathbf{2}_{\text{Tb}} - \tau_{\text{fast}}$	2.7×10^{-3}			1.0×10^{-5}	14.6
$\mathbf{2}_{\text{Tb}} - \tau_{\text{slow}}$		1.06	0.7	2.6×10^{-6}	23.6
$\mathbf{4}_{\text{Ho}} - \tau_{\text{slow}}$	0.1			3.1×10^{-7}	25.8



elevated temperatures magnetic relaxation happens predominantly via Orbach relaxation with effective barriers of $U_{\text{eff}} = 14.6$ K, 23.6 K (fast and slow process 2_{Tb}), and 25.8 K (4_{Ho}), for the three sets of relaxation processes shown in Figure 8. These values are comparable to the barrier heights observed in other butterfly-shaped heterometallic 3d–4f complexes.^[21]

The absence of a χ'' signal in 1_{V} directly indicates that slow magnetic relaxation in the V_2Ln_2 compounds can only be observed when paramagnetic Ln^{III} ions are present. Hence, a number of conclusions can be drawn from the presented ac susceptibility data: (i) Slow magnetic relaxation in zero field occurs only in 2_{Tb} , 3_{Dy} and 4_{Ho} , implying that the ligand field around the Ln^{III} ions is more suitable for oblate charge distributions. (ii) Despite being predominantly prolate, the 4f-orbitals of Tm^{III} in 6_{Tm} experience a considerable axial crystal field anisotropy, such that only small external magnetic fields are needed to suppress quantum tunnelling. (iii) Sizable magnetic coupling to the V^{III} moments is not sufficient to slow down the magnetic relaxation to the Hz or even sub-Hz scale.

Discussion

The discussion of 3d–4f coupling beyond dipole-dipole-interaction in the literature was initially based on a model proposed for a Gd_2Cu_4 complex by Kahn *et al.*^[49] with similar ideas being reported by Gatteschi *et al.*^[57–58] In summary, the 3d–4f superexchange interaction arises from two distinct exchange mechanisms. The first of these is electron transfer from the occupied 3d orbitals of the transition metal ion into the empty 5d orbitals of the lanthanide favouring a ferromagnetic spin alignment due to enhanced Hund's coupling in the charge transfer state, as long as the 4f shell is at least half-filled.^[49] The second is excitation of the transition metal 3d electrons into the partially filled lanthanide 4f orbitals leading to either a ferromagnetic or an antiferromagnetic contribution to J_{3d-4f} depending on whether the corresponding orbitals are orthogonal to each other or overlap. Taking into account the influence of single ion anisotropy on the 3d orbitals, this model was further developed to include also other transition metal ions with $S > 1/2$.^[6,14–15] Considering the magnitude of $J_{\text{V-Ln}}$ obtained from HF-EPR (Table 2), it becomes evident that dipolar interactions are negligible. Hence, the observation of antiferromagnetic interaction implies that electron exchange between 3d and 4f orbitals is the dominant coupling mechanism. This finding can be rationalised by the predominant π -character of the occupied V^{III} t_{2g} orbitals, which renders the ferromagnetic σ -type 3d–5f charge transfer rather inefficient^[59] and favours direct orbital overlap with the Ln^{III} 4f orbitals.^[60]

Figure 9(a) summarises the results of our HF-EPR studies by showing the dependence of $J_{\text{V-Ln}}$ and θ on the lanthanide ion. Of the studied complexes 2_{Tb} exhibits the largest magnetic interaction between V^{III} and Ln^{III} moments. This coupling becomes significantly weaker when Tb^{III} is replaced by Ho^{III} and subsequently increases for the heavier lanthanides Er^{III} and Tm^{III} . Interestingly, the relative angle between the vanadium and lanthanide main anisotropy axes θ exhibits the inverse trend.

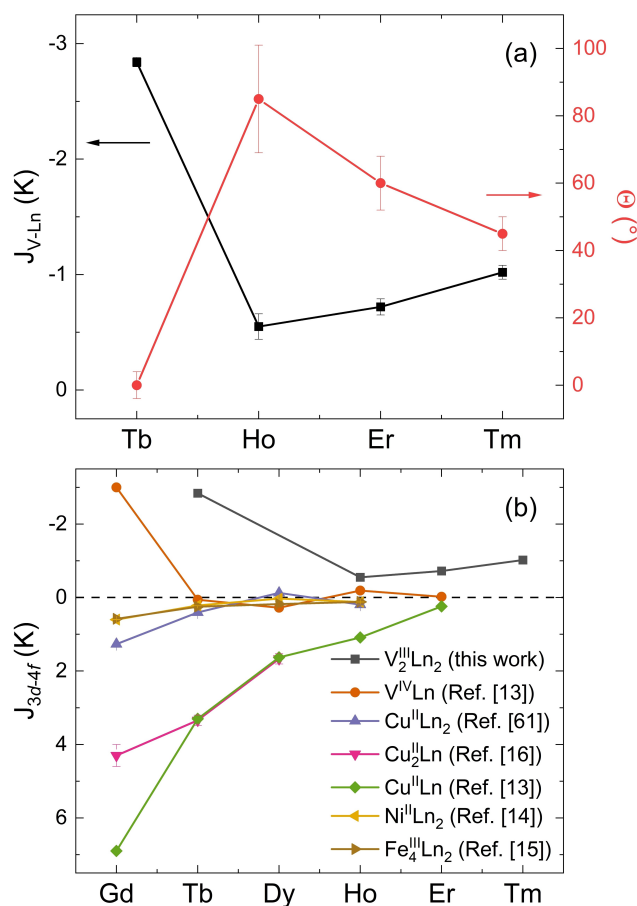


Figure 9. (a) Chemical trend of the $\text{V}^{\text{III}}\text{-Ln}^{\text{III}}$ coupling constant and of the relative angle between the main anisotropy axes (see equation 2) as determined by HF-EPR. (b) Experimentally determined J_{3d-4f} in different heterometallic complexes as reported in recent HF-EPR studies.

According to Kahn's model^[49] the ferromagnetic contribution to $J_{\text{V-Ln}}$ is expected to decrease with increasing atomic number from Gd^{III} onwards as the amount of unpaired 4f electrons, and hence Hund's coupling in the charge transfer state $4f^5 5d^1$, decreases along the series. While this mechanism qualitatively rationalises the increase of overall antiferromagnetic coupling from Ho^{III} to Tm^{III} , the finding of 2_{Tb} exhibiting the largest $J_{\text{V-Ln}}$ represents an outlier. The reason for this is (likely) to be the decreasing orbital overlap between the partially filled 3d and 4f orbitals due to a rotation of the Ln^{III} easy anisotropy axis (see Figure 9). A similar mechanism has been found to explain the dependence of $J_{\text{Cr-Gd}}$ on the Cr–F–Gd bond angle in $\text{Cr}^{\text{III}}\text{Gd}^{\text{III}}$ complexes, where the 3d–4f orbital overlap is mainly varied by a rotation of the Cr^{III} anisotropy tensor via changes in the ligand field.^[59]

In the literature, HF-EPR has established itself as the main method of choice for quantitative experimental determination of 3d–4f coupling.^[14–16,50–51,61–62] An overview of the $J_{\text{M-Ln}}$ (M = transition metal) in different heterometallic complexes as obtained in recent HF-EPR studies is given in Table 4 and visualised in Figure 9(b). In most of the investigated compounds, namely those with Cu^{II} , Ni^{II} and Fe^{III} as the 3d ion, ferromagnetic 3d–4f interaction is found and rationalised by a



Table 4. Magnetic 3d–4f coupling in selected heterometallic molecular complexes as reported here and in the literature.

Ln =	J_{M-Ln} (K)						
	Gd	Tb	Dy	Ho	Er	Tm	Ref.
$V^{III}_2Ln_2$		–2.84(4)		–0.55(11)	–0.72(7)	–1.02(6)	this work
$V^{IV}Ln$	–3.0	0.06(1)	0.28(1)	–0.19	–0.02(1)		13
$Cu^{II}Ln_2$	1.271(7)	0.405(3)	–0.126(3)	0.196(13)			61
Cu^{II}_2Ln	4.3(3)	3.34(14)	1.67(14)				16
$Cu^{II}Ln$	6.9	> 3.3	1.63(1)	1.09(2)	0.24(1)		13
$Ni^{II}Ln_2$	0.602(8)	0.216(12)	0.031	0.122(3)			14
$Fe^{III}_4Ln_2$	0.58	0.25(5)	0.18(8)	0.12(8)			15

dominance of the 3d to 5d charge transfer process.^[13–16,61] In contrast, the results on V^{III} -based butterfly complexes presented here evidence antiferromagnetic coupling between the 3d and 4f moments, which has also been observed in dinuclear V^{IV} -Ln molecules.^[13] As mentioned above, the qualitative differences probably arise from different spatial distributions of the contributing 3d orbitals, *e.g.*, $3d_{x^2-y^2}$ for Cu^{II} and $3d_{xy}$ for $V^{IV/III}$, that determine the absence or presence of overlap with the Ln^{III} orbitals.^[59–60] Similarly, the size of J_{M-Ln} depends on the coordination geometry and on the influence of other magnetic orbitals in the corresponding electronic configuration. Whereas the expected chemical trend of decreasing J_{M-Ln} for heavier lanthanide ions is generally confirmed by the experimental data, the studies are usually restricted to the series $Ln^{III} = Ce^{III}$ – Ho^{III} . Here, we expand this range to include the Ln^{III} ions beyond Gd^{III} with prolate electron density distribution, *i.e.* Er^{III} and Tm^{III} . In fact, the two regimes observed in Figure 9 can be related to the 4f anisotropy shape for the later Ln^{III} ions: For the oblate Ln ions J_{V-Ln} weakens with the Ln atomic number, but increases upon the series when the 4f electron distribution exhibits a prolate shape. The correlation of 4f anisotropy shape and magnetic 3d–4f coupling in the $V^{III}_2Ln_2$ butterfly series can be rationalised by considering changes in the Ln^{III} anisotropy axis. This leads to variation in orbital overlap as a result of ligand field effects, as described earlier.

Conclusions

We report the successful synthesis of a series of butterfly-motif $V^{III}_2Ln^{III}_2$ ($Ln^{III} = Tb^{III}$ – Yb^{III} , Y^{III}) complexes and their magnetic characterisation *via* combined magnetic SQUID, μ SQUID and HF-EPR studies. In particular, our analysis allows us to quantify 3d–4f coupling by treating the Ln^{III} moments as essentially Ising spins. Unlike for other 3d–4f complexes, J_{V-Ln} does not monotonically decrease upon increasing Ln atomic number but is shown to exhibit opposite effects for oblate and prolate Ln 4f orbital distributions. This trend is rationalised through the variation in orbital overlap due to a rotation of the Ln^{III} anisotropy axis. The important role of 4f anisotropy shape is also confirmed by the ac susceptibility data which show the absence of slow magnetic relaxation in zero field for prolate Ln^{III} ions.

Supporting Information

The authors have referred to reference^[53] and reference^[39] and cited an additional reference within the Supporting Information.^[63]

Details of the syntheses and crystallography together with Bond Valence Sum analyses and additional magnetic and HF-EPR data are provided in the Supplementary Information. Full crystallographic data and details of the structural determinations for the structures in this paper have been deposited with the Cambridge Crystallographic Data Centre as supplementary publication nos. CCDC 2406392–2406398. Copies of the data can be obtained, free of charge, from <https://www.ccdc.cam.ac.uk/structures/>

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Conflict of Interests

The authors declare no conflict of interest.

Data Availability Statement

Full crystallographic data and details of the structural determinations for the structures in this paper have been deposited with the Cambridge Crystallographic Data Centre as supplementary publication nos. CCDC 2406392–2406398. Copies of the data can be obtained, free of charge, from <https://www.ccdc.ca->



m.ac.uk/structures/. Magnetic data is available on request from the corresponding author.

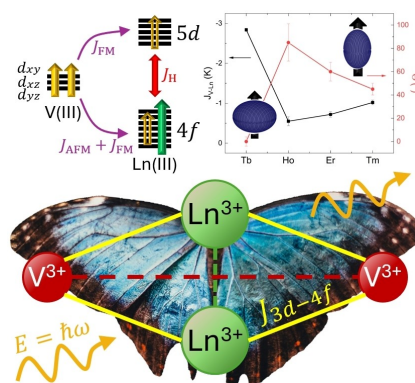
Keywords: 3d–4f Compounds • Exchange Interactions • HF-EPR • Magnetic Properties • Relaxation Mechanisms

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The 3d–4f coupling constants and relative angles between distinct anisotropy axes in butterfly-shaped $V^{III}_2Ln^{III}_2$ (Ln=Y, Tb–Yb) metal-organic molecular complexes are determined using combined high-field EPR spectroscopy and magnetisation studies. The observed chemical trend sheds light on the electronic exchange mechanism and highlights the important role of the 4f charge distribution.



J. Arneth, Dr. X.-F. Li, Dr. J. Braun*, Dr. S. Paul, Dr. M. Schulze, Dr. C. E. Anson, Prof. Dr. W. Wernsdorfer, Prof. Dr. A. K. Powell*, Prof. Dr. R. Klingeler**

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The Interplay of Single Ion Anisotropy and Magnetic 3d–4f Interactions in $V^{III}_2Ln^{III}_2$ Butterfly Complexes

